E. Hempel M. Beiner S. Höring E. Donth

# Struik law for enthalpy retardation at the glass transition in poly(*n*-alkylmethacrylates) measured by DSC aging experiments

Received: 26 January 1995 Accepted: 4 August 1995

Dr. E. Hempel (⋈) · M. Beiner · E. Donth Fachbereich Physik Universität Halle 06099 Halle/Saale, Germany

S. Höring Fachbereich Chemie Universität Halle 06099 Halle/Saale, Germany Abstract Results of calorimetric (DSC) experiments on a series of poly (n-alkylmethacrylates), from methyl to pentyl, after different aging times  $t_e$  at different aging temperatures  $T_e$  are presented. The aging behavior is quite different from that in other polymers, for example PS. For all poly(n-alkylmethacrylates) investigated the aging peak temperature  $T_{max}$  is shifted parallel to

the aging temperature  $T_{\rm e}$  in a large temperature interval below the glass temperature. The results are discussed with respect to shear and entropy response in the  $\alpha\beta$  splitting region.

**Key words** DSC – Struik law – physical aging –  $\alpha\beta$  splitting region – poly(*n*-alkylmethacrylates)

### Introduction

Poly(*n*-alkylmethacrylates) are of interest because the  $\alpha\beta$  splitting region, where  $\alpha$  and  $\beta$  traces are close together in the Arrhenius diagram (Fig. 1), is at low frequencies that are accessible for dynamic calorimetric and shear experiments [1, 2].

The position of the  $\alpha$  and  $\beta$  traces in the Arrhenius diagram of poly(ethylmethacrylate) (PEMA), poly(n-propylmethacrylate) (PnPrMA) and poly(n-butylmethacrylate) (PnBMA) was the subject of recent experimental work using photon correlation spectroscopy [3], dielectric and dynamic shear methods [4–7]. The  $\beta$  traces (local relaxation) of the poly(n-alkylmethacrylates) are quite similar, whereas the  $\alpha$  traces (cooperative relaxation) shift towards lower frequencies with increasing number of the C atoms in the side chain [8]. The shear compliance  $J^*$  (Fig. 2B) and the shear retardation spectrum L (Fig. 2A) for PMMA show that the shear compliance in the frequency range of the splitting region ( $\log \omega/s^{-1} > 3$ ) is higher than for polystyrene, PS (see also [9]). This is a hint for intensified relaxation processes between  $\alpha$  and  $\beta$  relaxation.

Physical aging experiments for shear in PEMA and PnBMA shift the  $\alpha$  relaxation in an Arrhenius diagram to higher temperatures. This corresponds qualitatively, after temperature time superposition, to the Struik law [10] in the Andrade zone between  $\alpha$  and  $\beta$  for many polymers, e.g., PS far below the splitting, describing some parallelism between the time shift of response  $(t_{meas})$  and the aging time  $(t_{\rm e}), \mu_{\rm Struik} = d\log t_{\rm meas}/d\log t_{\rm e} \approx {\rm const} \ (\approx 1).$  The high shear activity in the splitting region of the poly(n-alkylmethacrylates) suggests the question of if a similar behavior as observed in shear can, for these polymers, also be observed by calorimetry. If the material is "enthalpy active," or because of  $\Delta H \approx T \Delta S$  for isobaric experiments in condensed matter, "entropy active" in the splitting region – which means that aging varies the entropic state there – then we expect, similarly to shear, that the peak (partly called the overshoot in DSC experiments) is shifted with the aging time  $t_{\rm e}$  and aging temperature,  $T_{\rm e}$ .

The aim of this paper is a systematic experimental study of the DSC peak shift due to aging. Calorimetric aging experiments are described in the literature for PMMA [11] and recently for PEMA [12]. In this paper we include PnPrMA, PnBMA and poly(*n*-pentylmethacrylate)

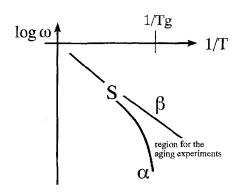
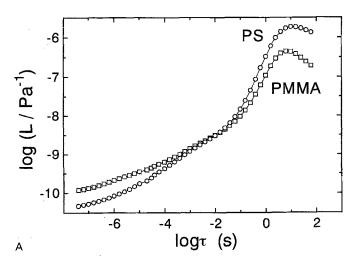
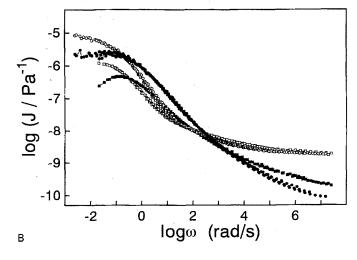


Fig. 1 Region for the aging experiments near the splitting region S of the Arrhenius diagram (schematically)

Fig. 2A Shear retardation spectrum L for PMMA (reduced to  $117\,^{\circ}\text{C}$ ) and PS (reduced to  $100\,^{\circ}\text{C}$ ), calculated by a nonlinear regularization method [13]. The symbol log always means the logarithm on a base 10,  $\log_{10}$  and  $\omega\tau = 1$ . B Mastered shear compliance J' and J'' data for PMMA (reduced to  $117\,^{\circ}\text{C}$ ) and PS (reduced to  $100\,^{\circ}\text{C}$ ) from which the spectra are calculated ( $\circ J'$ ,  $\bullet J''$  for PS;  $\Box J'$ ,  $\blacksquare J''$  for PMMA)





PnPMA. The results will be discussed in the framework of a modified Struik law, transferring  $\mu_{\text{Struik}}$  from the time dependence to a  $\tilde{\mu}$  for the temperature dependence.

## **Experimental**

The polymers PEMA, PnPrMA, PnBMA and the monomer *n*-pentylmethacrylate were purchased from Polyscience, Inc. They were prepared in the same way as described in [4] and [5]. PnPMA was obtained by radicalic polymerization. PS is a commercial product from Buna AG (formerly VEB Chemische Werke Buna)

The DSC glass transition temperatures  $T_{\rm g}$  (heating rate 10 K/min), the molecular weights  $M_{\rm w}$  and the  $M_{\rm w}/M_{\rm n}$  values (determined by GPC) for the poly(n-alkylmethacrylates) are listed in Table 1. The PS sample has a molecular weight of  $M_{\rm w}=360$  kg/mol. All poly(n-alkylmethacrylates) are atactic (78% syndio sequences in the average).

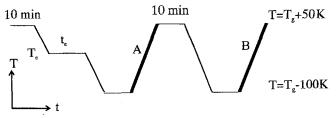
Calorimetric experiments were made in a DSC-7 instrument from Perkin Elmer. The sample weights were  $10 \dots 15$  mg, and the temperature time program is shown in Fig. 3. The cooling and heating rates were always 10 K/min. Dynamic shear compliance for PMMA and PS was measured in a RDA II instrument from Rheometrics Scientific using a stripe  $(1.5 \times 10 \times 25 \text{ mm}^3)$  as sample. The shear retardation spectrum was calculated using an algorithm according to [13]. The time dependent creep

Table I Molecular weight and glass transition temperature  $T_g$  for poly(n-alkylmethacrylates)

	$M_{\rm w}/{\rm kg\cdot mol^{-1}}$	$M_{ m w}/M_{ m n}$	$T_{\rm g}/{ m ^{\circ}C}$
PMMA	≈ 6700	1.8 ± 0.1	100 ± 1
PEMA	$72.3 \pm 3.0$	1.9 $\pm 0.1$	$74\pm1$
	$76 \pm 5$		
$PnPrMA^{a}$	$340 \pm 20$	$5.0 \pm 0.3$	$46 \pm 1$
PnBMA	$330 \pm 8$	$2.35 \pm 0.1$	$24 \pm 2$
PnPMA	$570 \pm 30$	$1.8 \pm 0.1$	$13 \pm 2$

a bimodal

Fig. 3 Temperature time program.  $T_{\rm e}$  aging temperature,  $t_{\rm e}$  aging time, A measuring run, B comparison run



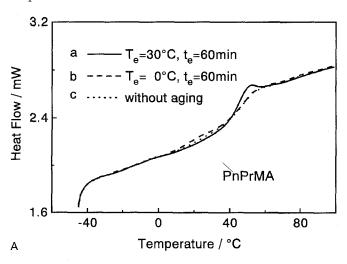
experiments on PEMA were performed in a DSR instrument from Rheometrics Scientific. The PEMA sample had a diameter of 8 mm and a thickness of about 3 mm.

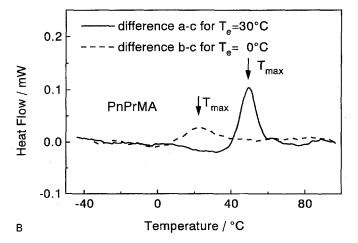
The polymers are measured as purchased or polymerized, respectively. The water content is small (<1%) and without detectable influence on the dynamic shear properties [4].

### Results

Figure 4A shows two examples for DSC runs after aging (curve a + b) with different aging temperatures  $T_e$ , and a comparison run (curve c) without aging for PnPrMA. The (a - c) and (b - c) difference curves have a peak with a maximum temperature  $T_{\text{max}}$  (Fig. 4B). In Fig. 5 the shift

Fig. 4A DSC curves (a, b) for two different aging temperatures  $T_{\rm e}$  and aging time  $t_{\rm e}=60$  min, and without aging (c) for PnPrMA. B Differences between the scans of Fig. 4A for the aged and unaged sample for PnPrMA





of  $T_{\rm max}$  is shown as function of the aging temperature  $T_{\rm e}$  for an aging time  $t_{\rm e}=60$  min. The accuracy of  $T_{\rm max}$  is about  $\pm 1$  K.

Normally [14–19] the shift of the overshoot with varying aging temperature is comparatively small (10 Kelvins), see for example the experimental points for PS in Fig. 5. But for the poly(n-alkylmethacrylates) the shift of  $T_{\rm max}$  is large. The relation between the maximum  $T_{\rm max}$  and the aging temperature  $T_{\rm e}$  (for a given aging time  $t_{\rm e} \ll t_{\rm equil}$ ) is similar for all the poly(n-alkylmethacrylates) in a temperature range  $0 < T_{\rm g} - T_{\rm e} < 50$  K (Fig. 5).

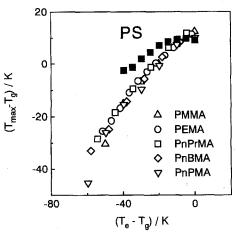
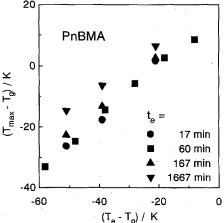


Fig. 5 Peak temperature  $T_{\rm max}$  as a function of aging temperature  $T_{\rm e}$  for an aging time  $t_{\rm e}=60$  min. The slopes  $\tilde{\mu}$  in the upper part of the diagram for PMMA, PEMA, PnPrMA, PnBMA and PnPMA are  $0.83\pm0.08$ ;  $0.76\pm0.04$ ;  $0.81\pm0.04$ ;  $0.85\pm0.03$  and  $0.88\pm0.09$ , respectively

Fig. 6 Peak temperature  $T_{\rm max}$  as a function of aging temperature  $T_{\rm e}$  for different aging times  $t_{\rm e}$  in PnBMA. The slopes  $\tilde{\mu}$  for  $t_{\rm e}=17$  min, 60 min, 167 min and 1667 min are  $0.90\pm0.09$ ;  $0.85\pm0.05$ ;  $0.84\pm0.05$  and  $0.70\pm0.05$ , respectively



The slope of the upper linear part,

$$\tilde{\mu} \equiv dT_{\text{max}}/dT_{\text{e}} = 0.85 \pm 0.10$$
  $(t_{\text{e}} = 60 \text{ min})$  (1)

is nearly the same for all the poly(n-alkylmethacrylates) investigated. The slope increases slightly for decreasing aging temperature  $T_{\rm e}$ . The shift of the annealing peaks with the aging time is shown in Fig. 6 for PnBMA. The slope  $\tilde{\mu}$  decreases with increasing annealing time  $t_{\rm e}$ . This finding confirms the result for PEMA given in [12]. The other poly(n-alkylmethacrylates) show a similar decrease of the slope with increasing annealing time.

# **Discussion and conclusion**

The reason for the  $T_{\rm max}$  shift is probably similar to that of the general shear-compliance shift reported by Struik [10], when translated from the time to the temperature variable by means of temperature time superposition. The compliance measurements shown in Fig. 2A + B indicate that the Andrade part in the compliance spectrum (slope of 0.3 for PMMA) corresponds to the frequency interval of the splitting region. In the temperature region between  $T_{\alpha}$  and  $T_{\beta}$  usually the Struik law [10] is valid ( $t_{\rm e}$  is the annealing time,  $t_{\rm meas}$  a characteristic relaxation time of the shear compliance,  $t_{\rm meas}(t_{\rm e})$ ):

$$\mu_{\text{Struik}} \equiv \frac{d \log t_{\text{meas}}}{d \log t_{\text{e}}} \approx 1 \ . \tag{2}$$

Equation (2) means that annealing time  $t_{\rm e}$  and the time  $t_{\rm meas}$  are scaled in the same way.

Comparing Eq. (1) and Eq. (2), the  $\mu$  ratio can be expressed by the ratio of the differential time-temperature (superposition) slopes

$$\frac{\mu_{\text{Struik}}}{\tilde{\mu}} = \left(\frac{d \log t_{\text{meas}}}{d T_{\text{max}}}\right) / \left(\frac{d \log t_{\text{e}}}{d T_{\text{e}}}\right). \tag{3}$$

According to Eq. 3, the equality  $\tilde{\mu} = \mu_{\text{Struik}}$  would mean that the temperature dependence of the (actual) relaxation time is equal to that of the annealing time, compared at a given aging state. Shear compliance experiments on PEMA (Fig. 7) result, however, in a somewhat different exponent,

$$\mu_{\text{Struik}} \approx 0.5 \pm 0.1$$
, (4)

i.e., for PEMA at 62 °C we have  $\mu_{\text{Struik}}/\tilde{\mu} \approx 0.66 \pm 0.11 < 1$ . This means that the actual and the aging time-temperature slope are not the same. This will be explained in Fig. 8 showing the equilibrium Arrhenius diagram for two selected processes (modes I, II) between  $\beta$  and  $\alpha$  in the splitting region. It is assumed that such processes can be defined (e.g., by sequential aging [20, 21]), and that there

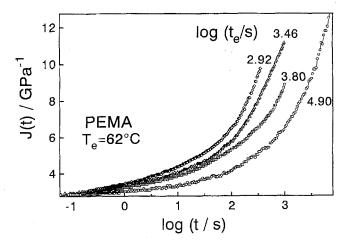


Fig. 7 Shift of shear compliance during physical aging for PEMA at  $T_{\rm e}=62\,^{\circ}{\rm C}$ . The parameter is the aging time,  $t_{\rm e}$ . The curves are slightly shifted in vertical direction to correct some rippling [4] between  $\alpha$  and  $\beta$ 

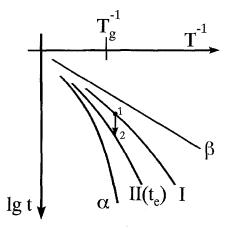


Fig. 8 Annealing  $1 \rightarrow 2$  in the splitting region between  $\alpha$  and  $\beta$ -traces [22] (schematically). The details are explained in the text

is a certain continuity of the slopes between  $\alpha$  and  $\beta$  (i.e., the curves are not parallel). Point 1 in Fig. 8 means [22] an actual state ( $T_{\rm meas}$ ,  $\lg t_{\rm meas}$ ), see also Eq. (3). Point 2 means a state reached after isothermal aging ( $T_{\rm e} = T_{\rm meas}$ ,  $t_{\rm e} > t_{\rm meas}$ ). We see that the slope dt/dT — mediating the time-temperature superposition — is really larger in point 2 than in point 1, which means, according to Eq. (3),  $\mu_{\rm Struik}/\tilde{\mu} < 1$ , as observed (under the condition that shear and entropy relaxes with the same rate).

It was also tried to reproduce the DSC traces by simulation with a Narayanaswamy program. Strong peak shifts can be simulated because of the small slopes of the traces in the splitting region of Fig. 8 [23]. But it is still an open question if the details of the DSC curves in the splitting region can be reproduced by such a program.

In summary, DSC aging experiments in the splitting region of poly(n-alkylmethacrylates) show pronounced peak shifts that can be explained by a temperature version of the Struik law originally formulated for shear in the time domain. Such an entropy activity indicates that molecular order observed by exchange NMR [24, 25] is also impor-

tant for calorimetry in the splitting region of these polymers.

Acknowledgements This work was partly supported by the Deutsche Forschungsgemeinschaft DFG. The authors thank Mr. S. Zeep for carrying out the polymerization and Mr. S. Reißig for calculation of the retardation spectra.

### References

- Heijboer J (1965) In: Prins JA. (ed) Physics of Non-Crystalline Solids (North-Holland, Amsterdam), p. 231
- McCrum NG, Read BE, Williams W (1991) Inelastic and Dielectric Effects in Polymeric Solids. Dover Publications, Inc., New York
- 3. Patterson GD, Jue PK, Ramsay DJ, Stevens JR (1994) J Polym Sci, Polym Phys Ed 32:1137-1147
- 4. Beiner M, Garwe F, Schröter K, Donth E (1994) Polymer 35:4127–4132
- Beiner M, Garwe F, Hempel E, Schawe J, Schröter K, Schönhals A, Donth E (1993) Physica A 201:72-78
- Garwe F, Beiner M, Hempel E, Schawe J, Schröter K, Schönhals A, Donth E (1994) J Non-Cryst Solids 172–174:191–196
- Garwe F, Schönhals A, Lockwenz H, Beiner M, Schröter K, Donth E (to be published)

- 8. Beiner M, Garwe F, Schröter K, Donth E (1994) J Colloid Polym Sci 272: 1439-1446
- 9. Plazek DJ, Craig AB, Chay IC (1994) J Non-Cryst Solids 172:181-190
- Struik LCE (1978) Physical Aging in Amorphous Polymers and Other Materials. Elsevier, Amsterdam, p 21
- Diaz-Calleja R, Perez J, Gomez-Ribelles JL (1989) Macromol Chem Macromol Symp 27:289-297
- 12. Estellés JM, Ribelles JLG, Prades MM (1993) Polymer 34:3837–3842
- 13. Honerkamp J, Weese J (1993) Rheol Acta 32:65-73
- 14. Hodge IM (1994) J Non-Cryst Solids 169:211-266
- 15. Wolkenstein MW, Sharonow JA (1961) Vysokomol Soedin 3:1739–1745
- Wunderlich B, Bodily DM, Kaplan MH (1964) J Appl Phys 35:95-101

- 17. Wolpert SM, Weitz A, Wunderlich B (1971) J Polym Sci [A-2] 9:1887
- 18. Illers K-H (1969) Macromol Chem 127:1-33
- De Bolt MA, Easteal AJ, Macedo PB, Moynihan CT (1976) J Amer Ceram Soc 59:16-21
- 20. McCrum NG (1984) Polymer Commun 25:2-4
- 21. Schönhals A, Donth E (1986) Acta Polym 37:475–480
- Donth E (1992) Relaxation and Thermodynamics in Polymers. Glass Transition. Akademie-Verlag, Berlin, p 196
- 23. Garwe F (1995) personal communica-
- Kulik AS, Beckham HW, Schmidt-Rohr K, Radloff D, Pawelzik U, Boeffel C, Spiess HW (1994) Macromolecules 27:4746-4754
- 25. Kulik AS, Radloff D, Spiess HW (1994) Macromolecules 27:3111-3113